

Designing Two-Qubit Gates for Exchange-Only Quantum Computation



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Work done in collaboration with:

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Zeuch, NEB, Phys. Rev. A **98**, 010303 (2016) Zeuch, NEB, Phys. Rev. B **102,** 075311 (2020)

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Early Vision of a Solid State Quantum Computer

Loss & DiVincenzo, Phys. Rev. B (1998)



Decades of Slow Steady Progress



Medford et al., Nature Nanotechnology (2013)



1 µm

1 µm

Andrews et al., Nature Nanotechnology (2019)



Basic Idea

• Use electron spins as qubits





spin-1/2 chain: electrons in quantum dots

• Quantum gates through spin exchange

$$H_i = J \; \boldsymbol{S}_i \cdot \boldsymbol{S}_{i+1}$$





spin-1/2 chain: electrons in quantum dots

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• Quantum gates through spin exchange

$$H_i = J \; \boldsymbol{S}_i \cdot \boldsymbol{S}_{i+1}$$





total spin of the oval

• Quantum gates through spin exchange

$$H_i = J \; \boldsymbol{S}_i \cdot \boldsymbol{S}_{i+1}$$





$$s_1 \otimes s_2 = |s_1 - s_2|, |s_1 - s_2 + 1|, ..., s_1 + s_2$$

Controlling Exchange

Petta et al., Science (2005)







Electron wave functions in quantum dot potential V(x)

Controlling Exchange



Electron wave functions in quantum dot potential V(x)

Controlling Exchange



Simple Exchange Pulses



• SWAP pulse $a = 0 \quad 1$ $t = \pi$ $n \quad t = \pi$ $\pi \rightarrow 0 \quad 1$ $1 \quad t = \pi$

Simple Exchange Pulses



Simple Exchange Pulses



- SWAP^{1/2} pulse

$$t = \pi/2$$

Three-Spin Qubit Encoding







DiVincenzo et al., Nature (2000)

Medford et al., Nature Nanotechnology (2013)

Three-Spin Qubit Encoding







DiVincenzo et al., Nature (2000)

Medford et al., Nature Nanotechnology (2013)

Single-Qubit Gates







DiVincenzo, Bacon, Kempe, Burkard & Whaley, Nature (2000)

19 pulse sequence found numerically



 $t_1 = 2.581..\ ,\ t_2 = 1.303..\ ,\ t_3 = 1.753..\ ,\ \ldots$

DiVincenzo, Bacon, Kempe, Burkard & Whaley, Nature (2000)

19 pulse sequence found numerically



 $t_1 = 2.581..$, $t_2 = 1.303..$, $t_3 = 1.753..$, ...

DiVincenzo, Bacon, Kempe, Burkard & Whaley, Nature (2000)



DiVincenzo, Bacon, Kempe, Burkard & Whaley, Nature (2000)



Fong-Wandzura Sequence

Fong & Wandzura, Quantum Information and Computation (2011)



Fong-Wandzura Sequence

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Two Simple Pulses


















Sequence Identity



Sequence Identity



Can show: Any pulse with $m^2=1$ will satisfy this identity

Sequence Identity



Can show: Any pulse with $m^2=1$ will satisfy this identity

Does *m* have to be a number? How about a matrix?

$m^2 = 1$ Pulses



 $m^2 = 1 \longrightarrow m = +1,-1$

$m^2 = 1$ Pulses



 $m^2 = 1 \longrightarrow m = +1,-1$

"Elevating" $m^2 = 1$ Pulses



"Elevating" $m^2 = 1$ Pulses



"Elevating" $m^2 = 1$ Pulses





 2×2 matrices, acting on "swapped in" qubit

 $M^2 = I$ "Pulse"



$M^2 = I \rightarrow M = +I, -I$

 $M^2 = I$ "Pulse"



$$M^{2} = I \rightarrow M = +I, -I \qquad M = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

More solutions! $\longrightarrow M = \hat{\boldsymbol{n}} \cdot \boldsymbol{\sigma} \qquad M = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \dots$











Applies *M* to bottom qubit only if top qubit is in state $|1\rangle$



Two-qubit controlled-*M* gate --- no leakage.



What we want









This is close! (But wrong basis)







Constraint



Any *V* satisfying this constraint will do the job.









For example, $t_1 = \pi$ and $t_2 = 0$



For example, $t_1 = \pi$ and $t_2 = 0$





For example, $t_1 = \pi$ and $t_2 = 0$





For example, $t_1 = \pi$ and $t_2 = 0$



$$\left\langle \left(\begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ \end{array} \right| \left| \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \\ 1 \end{array} \right| \left\langle \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ 1 \end{array} \right| \right\rangle = \left\langle \left(\begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ 1 \end{array} \right| \left| \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ \bullet \\ 1 \end{array} \right| \right\rangle \frac{1}{2} \left(-1 + e^{it_1} + e^{it_2} + e^{i(t_1 + t_2)} \right)$$

Need destructive interference

$$t_{1} = \frac{\pi}{2}$$

$$t_{2} = \frac{3\pi}{2}$$

$$e^{it_{1}} = i$$

$$e^{i(t_{1}+t_{2})} = 1$$

$$e^{it_{2}} = i$$

$$e^{it_{2}} = -i$$



Constraint



Optimal Solution







Constraint



Optimal Solution



Full Sequence



Full Sequence



Full Sequence



• <u>Original version</u> of the Fong-Wandzura Sequence



PHYSICAL REVIEW A 93, 010303(R) (2016)

Simple derivation of the Fong-Wandzura pulse sequence

Daniel Zeuch and N. E. Bonesteel

Department of Physics and National High Magnetic Field Laboratory, Florida State University, Tallahassee, Florida 32310, USA (Received 4 September 2015; published 25 January 2016)

We give an analytic construction of a class of two-qubit gate pulse sequences that act on five of the six spin- $\frac{1}{2}$ particles used to encode a pair of exchange-only three-spin qubits. Within this class, the problem of gate construction reduces to that of finding a smaller sequence that acts on four spins and is subject to a simple constraint. The optimal sequence satisfying this constraint yields a two-qubit gate sequence equivalent to that found numerically by Fong and Wandzura. Our construction is sufficiently simple that it can be carried out entirely with pen, paper, and knowledge of a few basic facts about quantum spin.

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Efficient two-qubit pulse sequences beyond CNOT

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We design efficient controlled-rotation gates with arbitrary angle acting on three-spin encoded qubits for exchange-only quantum computation. Two pulse sequence constructions are given. The first is motivated by an analytic derivation of the efficient Fong-Wandzura sequence for an exact CNOT gate. This derivation, briefly reviewed here, is based on elevating short sequences of SWAP pulses to an entangling two-qubit gate. To go beyond CNOT, we apply a similar elevation to a modified short sequence consisting of SWAPs and one pulse of arbitrary duration. This results in two-qubit sequences that carry out controlled-rotation gates of arbitrary angle. The second construction streamlines a class of arbitrary CPHASE gates established earlier. Both constructions are based on building two-qubit sequences out of subsequences with special properties that render each step of the construction analytically tractable.

DOI: 10.1103/PhysRevB.102.075311

New Sequences?

Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.

Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.



Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.



Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.



Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.



Sequence Identity

Strategy: "Elevate" simple three-spin pulse sequence to a two-qubit gate.



"Elevating" t Pulse



 $\begin{pmatrix} 1 \\ e^{it} \end{pmatrix}$

M can be any 2x2 unitary. No longer require $M^2 = I$



"Elevating" SWAP Pulse









Applies *M* to bottom qubit only if top qubit is in state $|1\rangle$





Constructing T



$$a = 0 \quad 1$$

$$M(t) = e^{i\xi(t)}e^{i\phi(t)\hat{n}(t)\cdot\sigma/2}$$

$$\phi(t) = 2 \arccos((5\cos(t/2) + 3\cos(3t/2))/8)$$

$$\xi(t) = -t/2$$

Constructing Π



Constraint



Solution



Constructing Π



$$t_1$$
=1.34004..., s_1 = 2 π - $t_{1,...}$



Constraint



Solution



Full Sequence



Full Sequence



 $t_1 = 1.34004..., s_1 = 2\pi - t_{1,...}$

 $\phi(t) = 2 \arccos((5 \cos(t/2) + 3 \cos(3t/2))/8)$

Summary

Analytic Derivation of Fong-Wandzura CNOT sequence:

Open questions:

- 1) Can we prove Fong-Wandzura sequence is truly optimal?
- 2) More efficient general gate constructions?
- 3) Can these tools be used to construct more "robust" sequences?

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